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Molecular interaction studies [and](http://www.elsevier.com/locate/tca) [theoretical](http://www.elsevier.com/locate/tca) [estima](http://www.elsevier.com/locate/tca)tion of ultrasonic speeds using scaled particle theory in binary mixtures of toluene with homologous nitriles at different temperatures

K. Rajagopal^a, S. Chenthilnath^{b,∗}

^a Department of Physics, Govt. College of Engineering, Tirunelveli 627007, Tamilnadu, India

^b Department of Physics, Satyam College of Engineering and Technology, Nagercoil 629301, Tamilnadu, India

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ABSTRACT

Densities, ρ , ultrasonic speeds, u, and viscosities, η of binary mixtures of toluene with acetonitrile (AN), propionitrile (PN) and butyronitrile (BN), including those of pure liquids are measured over the entire composition range at temperatures 298.15, 303.15 and 308.15 K, respectively. From these experimental data, excess intermolecular free length, L_f^E , the deviations in ultrasonic speed, Δu , excess acoustic impedance, Z^E and excess free energy of activation ΔG^{*E} are calculated. The variation of these excess parameters indicates the presence of specific interactions between toluene and AN/PN/BN molecules, and the interactions follow the order: AN > PN > BN, i.e., the toluene–nitrile interaction decreases with increase in alkyl chain length. Moreover, the theoretical ultrasonic speeds are computed based on the scaled particle theory and compared with the experimentally measured values.

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1. Introduction

Molecular properties in liquid state are very useful for design calculations involving separations, heat transfer, mass transfer and fluid flow. Mixing volume effects play a pivotal role in many practical applications that include paints, varnishes and printing ink industries [1,2]. The present work is in continuation of our earlier studies [3,4] on understanding thermodynamic properties of binary mixtures whose components have relevant industrial applications.

Literature survey shows that molecular interaction of homologous series of nitriles with toluene using volumetric, viscometric [and](#page-8-0) [a](#page-8-0)coustic studies have not been reported except for the con[tra](#page-8-0)dicting reports in AN + toluene binary mixture [5,6]. Abraham et al. [5] concluded the presence of strong interactions between AN–toluene molecules from negative deviations in isentropic compressibility studies. However, Ritzoulis et al. [6] reported contradicting weak interactions between the same AN–toluene binary mixture using excess viscosity values. [Thus](#page-8-0) [th](#page-8-0)e above contradicting results promoted us to reinvestigate the molecular interactions in the above system other than tolue[ne + P](#page-8-0)N/BN systems at three different temperatures.

Nitrile molecules are aprotic and highly polar with their dipoles oriented anti-parallel to each other and show strongly ordered structure due to strong dipole–dipole interactions [7]. The toluene molecules possess large quadrupole moments, causing an orientational ordering in this liquid. The orientational order is thought of as a partial alignment of neighbouring segments or possibly of whole molecules [8]. The binary mixtures containing aromatic hydrocarbons are interesting because they find a[pplic](#page-8-0)ations in the studies of polymer phase diagrams and preferential interaction of polymers in mixed solvents [9,10]. Therefore, the study of intermolecular interactions in toluene + homologous aliphatic nitrile mixtures would [be](#page-8-0) [in](#page-8-0)teresting owing to their industrial applications.

In this paper, we report the densities, ρ , ultrasonic speeds, u and viscosities, η of pure toluene, AN, PN and BN and those of thei[r](#page-8-0) [binary](#page-8-0) mixtures over the entire composition range at 298.15, 303.15 and 308.15 K. From these experimental data, the values of $L_f^{\tilde{E}}$, Δu , Z^E and ΔG^{*E} are calculated and discussed in terms of molecular interactions between the unlike molecules. Furthermore, theoretical ultrasonic speeds of thesemixtures are evaluated by considering the different shapes of the participating components using the scaled particle theory and compared with the experimental values.

2. Experimental

Toluene (Ranchem, India, purity >99.7%), AN, PN and BN (E. Merck, India, purity >99%) were purified by using the methods

[∗] Corresponding author. Tel.: +91 9443 418697; fax: +91 4652 227621. E-mail address: chenthilnaths@rediffmail.com (S. Chenthilnath).

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Table 1 Comparison of experimental values of density, ρ , viscosity, η and ultrasonic speed, u of pure liquids with the corresponding literature values at different temperatures.

Liquids	T(K)	ρ (10 ³ kg m ⁻³)		η (10 ⁻³ N s m ⁻²)		$u (m s^{-1})$	
		Exp	Lit	Exp	Lit	Exp	Lit
Toluene	298.15	0.8609	0.862055 [15]	0.5635	0.554 [16]	1299.8	1304 [17-19]
	303.15	0.8560	0.8566 [20]	0.5344	0.5284 [21]	1279.5	1282 [21]
	308.15	0.8516	0.85184 [20]	0.5091	0.5068 [22]	1260.7	1259 [23]
AN	298.15	0.7753	0.7764 [24]	0.3696	0.341 [25]	1271.3	1268.5 [25]
	303.15	0.7699	0.77122 [26]	0.3554	0.3501 [27]	1251.9	1258.64 [26]
	308.15	0.7647	0.7662 [28]	0.3414	0.3285 [5]	1231.7	1239 [5,29]
PN	298.15	0.7754	$-$	0.4107		1252.4	
	303.15	0.7699	0.77191 [26]	0.3991	0.389 [12]	1232.7	1237.63 [26]
	308.15	0.7652	0.7646 [30]	0.3440		1213.7	1223.7 [30]
BN	298.15	0.7847		0.5463		1271.4	
	303.15	0.7803	0.78172 [26]	0.5255	0.515 [12]	1252.2	1260 [31]
	308.15	0.7759	$\qquad \qquad -$	0.5049	$\qquad \qquad -$	1233.9	$\qquad \qquad -$

described in the literature [11,12] and only middle fractions were used in the experiment. All the chemicals were st[ored](#page-8-0) over 0.4 nm molecular sieves to remove water content and degassed just before use. The mixtures were prepared by mass and were kept in special airtight stopper glass bottles to avoid evaporation. The weighings were done by [using](#page-8-0) [an](#page-8-0) electronic balance (AND, HR series 300, Japan) with a precision of ± 0.1 mg. The uncertainty in the mole fraction was estimated to be less than $\pm 1 \times 10^{-4}$.

Densities of the pure liquids and their binary mixtures were measured by using a single capillary pycnometer (Borosil glass) having a bulb capacity of 10 mL. The capillary with graduated stem marks had a uniform bore and could be closed by a well-fitted glass cap. The marks on the capillary were calibrated with triple distilled water. The reproducibility in the density measurements was within \pm 0.1 kg m⁻³ [13]. The ultrasonic speeds in pure liquids and in their binary mixtures were measured using a single-crystal variable-path multifrequency ultrasonic interferometer operated at 2 MHz. The reproducibility in the ultrasonic speedmeasurements was within ± 0.03 %. The viscosities of pure liquids and their binary mixtures we[re](#page-8-0) [mea](#page-8-0)sured by using Ubbelohde type suspended level viscometer with flow time of 174 s for triply distilled water at 308.15 K. As the flow time was greater than 100 s, the kinetic energy corrections were not necessary [14]. The viscometer containing the test liquid was allowed to stand for about 30 min in a thermostatic water bath so that the thermal fluctuations in viscometer were minimized. The time of flow was measured with an electronic stopwatch capable of recording ± 0.01 s. An average of four sets of flow times for each mix[ture](#page-8-0) [w](#page-8-0)as taken for the purpose of calculation of the viscosity. The overall experimental reproducibility was estimated to be $\pm 2.0 \times 10^{-3}$ Ns m⁻² [13]. The temperature of the test liquids during the measurements was maintained within an uncertainty of \pm 0.01 K in an electronically controlled thermostatic water bath.

The reliability of experimental measurements of ρ , u and η were ascertained by comp[aring](#page-8-0) [t](#page-8-0)he experimental data of pure liquids with the corresponding literature values (see Table 1).

3. Theory

The thermodynamic parameters like intermolecular free length (L_f) , acoustic impedance (Z) and free energy of activation (ΔG^*) are calculated using the following standard relations [32,33]:

$$
L_f = \frac{K}{u\rho^{1/2}}\tag{1}
$$

$$
Z = u\rho \tag{2}
$$

$$
\Delta G^* = RT \ln \left[\frac{\eta V}{hN} \right] \tag{3}
$$

where K is Jacobson's constant that depends upon temperature and is given by $[K = (93.875 + 0.375T) \times 10^{-8}]$ $[K = (93.875 + 0.375T) \times 10^{-8}]$ $[K = (93.875 + 0.375T) \times 10^{-8}]$ [\[34\],](#page-8-0) R is the universal gas constant, T being the absolute temperature. V is the molar volume $[M_{\text{eff}}/\rho]$, M_{eff} is the effective molecular weight $[M_{eff} = (x₁M₁ + x₂M₂)]$, where $M₁$, $M₂$, $x₁$ and $x₂$ are the molecular weights and the mole fractions of toluene and nitriles, respectively. h and N are the Planck's const[ant](#page-8-0) [an](#page-8-0)d the Avogadro number, respectively. Instead of thermodynamic parameters, thermodynamic excess functions are found to be very sensitive towards mutual interactions between the component molecules of the binary mixtures. The sign and the extent of deviation of the functions from ideality depend on the strength of interactions between unlike molecules [35,36]. The excess intermolecular free length, L_f^E ,

excess acoustic impedance, Z^E and excess free energy of activation, ΔG^{*E} are calculated using the general relation:

$$
Y^{E} = Y - [x_1 Y_1 + x_2 Y_2]
$$
 (4a)

where *Y* represents L_f/Z or ΔG^* of the mixture.

The deviations in ultrasonic speed, Δu , may be estimated using

$$
\Delta u = u - [x_1 u_1 + x_2 u_2]
$$
 (4b)

The excess functions L_f^E , Δu , Z^E and ΔG^{*E} are fitted to a Redlich–Kister type [37] polynomial equation:

$$
P^{E} = x(1-x)\sum_{i=0}^{n} A_{i}(1-2x)^{i-1}
$$
\n(5)

where P^{[E](#page-8-0)} is $L_f^E/\Delta u/Z^E$ and ΔG^{*E} . The values of A_i are evaluated by using least squares method with all points weighed equally and the corresponding standard deviations $\sigma(P^E)$ are calculated by using the relation:

$$
\sigma(P^E) = \sum \left(\frac{\left(Y_{\text{cal}}^E - Y_{\text{exp}}^E\right)^2}{n - j} \right)^{1/2} \tag{6}
$$

where n is the number of experimental data points and j is the number of A_i coefficients considered. The values of P^E_{cal} are obtained from Eq. (5) by using the best-fit values of A_i coefficients.

4. Results

The experimental values of densities, ρ and ultrasonic speeds, u and viscosities, η of binary mixtures of toluene with AN, PN and BN, over the entire composition range, expressed in terms of mole fraction, x_1 of toluene at various temperatures are listed in Table 2. The values of coefficients A_0 , A_1 , A_2 , A_3 and A_4 along with standard deviation σ of fit for all the mixtures are listed in Table 3. The

Table 2

Densities, ρ , ultrasonic speeds, u , viscosities, η , deviations in ultrasonic speed, Δu , excess intermolecular free length, L_f^E , excess acoustic impedance, Z^E and excess free energy of activation, ΔG^{*E} as functions of mole fraction, x_1 of toluene for toluene + AN/PN/BN mixtures at different temperatures.

variations of L_f^E , Δu , Z^E and ΔG^{*E} with mole fraction, x_1 of toluene at 298.15 K, along with the smoothed values by using Eq. (5) are presented graphically in Figs. 1–4.

5. Discussion

5.1. Excess thermodynamic parameters

The increase in u values and decrease in L_f values with x_1 (see Table 2) indicates significant interaction between the mixtures [38]. From Figs. 1–3, it is seen that L_f^E values are negative over the entire mole fraction range and at all temperatures investigated while Δu and Z^E show opposite trends indicating the presence of specific interactions between toluene and nitrile molecules [39]. The magnitude of L_f^E values for these mixtures follow the order: AN < PN < BN. In general, the observed L_f^E , Δu and Z^E values depend on several factors, which are physical and chemical in nature [7,40]. The physical contributions, i.e., dispersion forces [and no](#page-8-0)n-specific physical (weak) interactions that lead to positive L_f^E and negative Δu and Z^E values. Further physical contributions are also due to geometrical effects which allow molecules with v[ery diffe](#page-8-0)rent sizes to fit into each other's structure resulting in negative L_f^E and positive Δu and Z^E values. Chemical contributions involve breaking up of the hydrogen-bonded structure(s), resulting in positive $L_f^{\bar{E}}$

Fig. 1. Plots of L_f^E against mole fraction, x_1 of toluene for binary mixtures of toluene with AN (\blacklozenge), PN (\blacksquare) and BN (\blacktriangle) at 298.15 K.

Table 3 Coefficients of A_i of Eq. (5) and standard deviations σ for the binary mixtures at different temperatures.

T/(K)	A_0	A ₁	A ₂	A_3	A_4	σ
Toluene + AN L_{ϵ}^{E} (10 ⁻¹¹ m)						
298.15	-0.3925	-0.1439	0.0610	0.0338	-0.1016	0.0002
303.15	-3.7522	-1.2575	0.1710	-0.0151	0.0885	0.0033
308.15	-3.7286	-1.4687	1.0031	0.8731	-1.1115	0.0017
Δu (m s ⁻¹) 298.15	2.3269					0.0040
303.15	1.8199	0.8987 0.8393	-1.8804 -1.3329	-0.8306 -0.8933	1.8423 0.5477	0.0056
308.15	1.5294	0.8992	-2.5448	-1.7581	2.2068	0.0036
	Z^{E} (10 ⁵ kg m ⁻² s ⁻¹)					
298.15	1.1675	0.3524	-0.1371	-0.0702	0.2093	0.0005
303.15 308.15	1.0746	0.2753	-0.0202 -0.2034	0.0665	-0.0717	0.0007
	1.0202	0.3222		-0.1550	0.2146	0.0003
	ΔG^{*E} (k[mol ⁻¹)					
298.15	1.0456	0.0873	0.0157	0.1337	-0.1308	0.0005
303.15	1.0081	0.0769	0.0181	0.1679	-0.2109	0.0002
308.15	0.9740	0.1201	-0.0589	0.0778	-0.1031	0.0006
Toluene + PN L_{ϵ}^{E} (10 ⁻¹¹ m)						
298.15	-3.3150	-0.9782	1.1273	0.5582	-1.1849	0.0020
303.15	-3.2302	-0.8166	1.1601	0.4350	-0.8746	0.0033
308.15	-3.1180	-0.7629	1.1399	0.4831	-0.5938	0.0008
Δu (m s ⁻¹) 298.15	2.9394	1.1156	-2.7312	-1.1301	2.7046	0.0043
303.15	2.5345	0.8433	-2.6032	-0.9310	2.0173	0.0064
308.15	2.0935	0.7558	-2.6086	-1.0289	1.8959	0.0030
	Z^{E} (10 ⁵ kg m ⁻² s ⁻¹)					
298.15 303.15	0.8080 0.7512	0.1657 0.1172	-0.2477 -0.2484	-0.1194 -0.0676	0.2386 0.1763	0.0003 0.0006
308.15	0.6942	0.0980	-0.2202	-0.0609	0.0819	0.0003
	ΔG^{*E} (k[mol ⁻¹)					
298.15	0.4248	0.0735	-0.0954	-0.0159	0.0817	0.0003
303.15	3.7469	0.6070	-0.7022	-0.0953	0.0666	0.0031
308.15	3.4076	0.4328	-0.6483	-0.1313	-0.0577	0.0026
Toluene + BN L_f^E (10 ⁻¹¹ m)						
298.15	-2.5489	-0.4781	1.0541	0.0577	-0.9840	0.0026
303.15	-2.5014	-0.2586	1.2341	-0.1142	-0.9045	0.0041
308.15	-2.3441	-0.4228	1.3262	0.0525	-1.0476	0.0035
Δu (m s ⁻¹)						
298.15	3.4050	0.7668	-2.3038	-0.0561	2.0896	0.0050
303.15	3.1478	0.2629	-2.4107	0.3504	1.7628	0.0089
308.15	2.6439	0.6897	-2.3216	-0.1708	1.8549	0.0070
	Z^{E} (10 ⁵ kg m ⁻² s ⁻¹)					
298.15	0.5741	0.0522	-0.2314	-0.0100	0.1999	0.0006
303.15	0.5368	0.0092	-0.2626	0.0276	0.1761	0.0004
308.15	0.4838	0.0349	-0.2861	0.0073	0.2106	0.0006
298.15	ΔG^{*E} (kJ mol ⁻¹) 0.0554	0.0094	-0.0036	0.0043	0.0175	0.0000
303.15	0.0460	0.0120	-0.0224	-0.0063	0.0173	0.0000
308.15	0.0331	0.0076	-0.0155	-0.0055	0.0097	0.0001

and negative Δu and Z^E values. However, specific interactions, such as the formation of hydrogen bonds, formation of charge-transfer complexes and strong dipole–dipole interactions between component molecules, result in negative L_f^E and positive Δu and Z^E values. Thus the observed negative trends in L_f^E and positive trends in Δu

and Z^E values accounts for strong specific interaction between the unlike molecules of the mixtures.

A plausible qualitative interpretation of the behaviour of these mixtures with composition has been suggested. As stated earlier, the AN molecules have strong dipole–dipole interactions while toluene molecules possess large quadrupole moments, causing an

Fig. 2. Plots of Δu against mole fraction, x_1 of toluene for binary mixtures of toluene with AN (♦), PN (■) and BN (▲) at 298.15 K

orientational order. Mixing of toluene with nitriles (AN/PN/BN) induces a decrease in the molecular order in toluene and dissociation of dipole–dipole association in nitriles, resulting in an expansion in volume, and hence, positive L_f^E and negative Δu and Z^E values. On the other hand, there is possibility of the electron donor–acceptor type (charge-transfer) interactions [41] between highly electronegative nitrogen of nitriles (acting as acceptor) and the π -electrons of ring of toluene molecules (acting as donor), resulting in negative $\overline{L_f^E}$ and positive Δu and Z^E values. Thus the observed negative \tilde{L}^E_f [and](#page-8-0) positive Δu and Z^E values suggest the presence of significant donor–acceptor interactions between toluene and AN/PN/BN molecules. Ma et al. [41,42] have also reported similar type of donor–acceptor interactions between elec-

Fig. 3. Plots of Z^E against mole fraction, x_1 of toluene for binary mixtures of toluene with AN (♦), PN (■) and BN (▲) at 298.15 K

Fig. 4. Plots of ΔG^{*E} against mole fraction, x_1 of toluene for binary mixtures of toluene with AN (\blacklozenge), PN (\blacksquare) and BN (\blacktriangle) at 298.15 K.

tronegative oxygen atom of sulpholane and π -electrons of the aromatic hydrocarbons.

Another contribution to negative L_f^E comes from the fitting of small AN/PN/BN molecules (molar volumes: $AN = 52.947 \text{ cm}^3 \text{ mol}^{-1}$, $PN = 71.034 \text{ cm}^3 \text{ mol}^{-1}$, BN = 88.072 cm³ mol⁻¹ at T = 298.15 K) into the voids created by bigger toluene molecules (molar volume = $107.03 \text{ cm}^3 \text{ mol}^{-1}$ at $T = 298.15$ K) [22]. Contributions arising from the fitting of smaller molecules into the voids available in the structure of bigger molecules were also considered by others [35,36,43] for interpreting excess parameter variations. Furthermore, the magnitude of L_f^E [value](#page-8-0)s at equimolar composition of these mixtures at 298.15 K follows the order: AN < PN < BN, which, in turn, indicates the order of the interactions between tol[uene and nit](#page-8-0)rile molecules.

It is interesting to note that the interaction between toluene and nitrile molecule in the mixtures decreases with increase in alkyl chain length of nitrile molecules. This may be due to the electron donor effect of alkyl group. Our observations are further supported by the results reported by Roy et al. [44] that large differences in dielectric constant of participating components of the mixture may result in negative L_f^E values. The dielectric constant of toluene is 2.379 at 296.15 K, while those of AN, PN and BN are 36.64, 29.70 and 24.83, respectively at 29[3.15 K](#page-8-0) [45]. The differences in dielectric constant values of toluene and nitrile molecules are large and may also result in negative L_f^E and positive Δu and Z^E values for the systems under study. Similar conclusions regarding the L_f^E , Δu and Z^E values were report[ed for](#page-8-0) N,N-dimethylformamide + 2methoxyethanol and N,N-dimethylformamide + 2-ethoxyethanol binary mixtures by Alisha et al. [46].

It has been reported earlier $[47]$ that, the viscosity and hence the free energy of activation of a mixture strongly depends on entropy of mixture, which is related with liquid's structure and enthalpy (and consequently with molecular interactions between the components of the mixt[ure\).](#page-8-0) [S](#page-8-0)o, the excess free energy of activation variations are fun[ctions](#page-8-0) of molecular interactions as well as size and shape of molecules.

According to Fort and Moore [35], excess free energy of activation and deviation in viscosity tends to become more positive as the strength of interaction increases. In the present investiga-

Table 4 Molecular assignment for different shapes.

Shape	Size	\overline{R}	S	V_H
Sphere Cube	Radius = a $Side = l$	\overline{a} 31/4	$4\pi a^2$ $6l^2$	$4\pi a^3/3$ 1^3
Tetrahedron	$Side = l$	$\left(3l\arctan\sqrt{2}\right)/2\pi$	$\sqrt{3}l^2$	$(\sqrt{2}/12)l^3$
Discs		radius = a and depth l		
Disc A	$l = a$	$(\pi + 1)a/4$	$4\pi a^2$	πa^3
Disc B	$l = a/4$	$(\pi + 0.25)a/4$	$5\pi a^2/2$	$\pi a^3/4$
Disc C	$l = a/2$	$(\pi + 0.50)a/4$	$3\pi a^2$	$\pi a^3/2$
Disc D	$l = a/10$	$(\pi + 0.10)a/4$	$11\pi a^2/5$	$\pi a^3/10$

tion, the positive trends in ΔG^{*E} (Fig. 4) indicate the presence of specific interactions between unlike molecules, thereby complimenting the results obtained through L_f^E , Δu and Z^E . Similar positive ΔG^{*E} values accounting strong interactions have been reported for DMSO + propionic acid and DMSO + n-butyric acid binary mixtures by Mishra et al. [33].

5.2. Scaled particle theory

The theoretical ultrasonic speeds of binary mixtures may be esti[mated](#page-8-0) based on some empirical, semi empirical and statistical models such as FLT, CFT, Nomoto and Vandeel Vangeal for the binary mixtures [48]. However, the common drawback of these theoretical models is that, the shapes of the participating species have not been taken into consideration. However, in scaled particle theory (SPT) [49] different shapes (such as sphere, cube, tetrahedral, disc A, disc B, disc C and disc D) of the participating components are c[onside](#page-8-0)red and when the participating components have the correct shapes the theoretical ultrasonic speed estimated based on this model will give values close to the experimental values. Gen[erall](#page-8-0)y, the chemical structure of a liquid molecule is known but no definite shape has been attached to it [50]. Recently, Ghosh et al. [50] have used the scaled particle theory for the binary mixtures of 1,1,1-trichloroethane with 1-alkanols by considering different shapes such as spherical, cubical and tetrahedral. However, Kalidoss et al. [51] also have used the SPT for some binary mixtures and ternary liquid mixtures by c[onside](#page-8-0)ring seven shapes of the individual components that includes sphere, cube, tetrahedron, disc A, disc B, disc C and disc D (the shapes are assigned by some characteristic parameters as given in Tables 4 and 5, respectively). In this paper, [theo](#page-8-0)retical ultrasonic speeds of the binary mixtures are estimated by assigning the above said seven shapes to the participating components and compared them with experimental speeds. Only when the participating molecules have particular shapes the theoretical values closely agrees with experimental ultrasonic speeds that will be determined by Chi-square fit (χ^2) [52].

The equation of state of fluid in the scaled particle theory is

$$
\frac{p}{\rho_N k_B T} = \frac{1 + \alpha + \alpha^2}{\left(1 - \alpha\right)^2} \tag{7}
$$

where α = $V_H \rho_N$, V_H is har[d](#page-8-0) [core](#page-8-0) volume, ρ_N [51] is number density and the other quantities have usual meaning [50].

Table 6 Values of a (by Eq. (13)).

298.15K				303.15K				308.15K			
TOL	AN	PN	BN	TOL	AN	PN	BN	TOL	AN	PN	BN
0.462	0.363	0.393	0.422	0.457	0.358	0.388	0.418	0.453	0.353	0.384	0.413
0.399	0.304	0.332	0.361	0.394	0.299	0.328	0.356	0.390	0.294	0.323	0.352
0.337	0.247	0.274	0.300	0.332	0.243	0.269	0.296	0.328	0.239	0.265	0.292
0.412	0.316	0.345	0.373	0.407	0.311	0.340	0.369	0.403	0.306	0.335	0.364
0.300	0.216	0.241	0.265	0.296	0.212	0.237	0.261	0.292	0.208	0.233	0.258
0.369	0.276	0.304	0.331	0.364	0.271	0.299	0.327	0.360	0.267	0.295	0.322
0.196	0.131	0.149	0.168	0.192	0.129	0.146	0.165	0.189	0.126	0.143	0.162

Table 7

Ultrasonic speeds computed theoretically using SPT with behavioural shape having least Chi-square fit and experimentally measured values of the binary mixture TOL + AN at 298.15 K.

x_1	$TOL+AN$											
	u_{\exp}	u_{cal}										
		Disc D + sphere	Disc D + cube	Disc D + tetra	Disc $D +$ disc A	Disc $D +$ disc B	Disc $D +$ disc C	Disc $D +$ disc D				
0.0000	1271.3	1271.5	1271.3	1271.2	1250.4	1271.3	1271.4	1271.2				
0.0506	1273.8	1265.8	1265.5	1267.7	1248.5	1265.0	1265.0	1265.6				
0.1082	1276.5	1265.7	1264.7	1267.6	1251.1	1263.4	1263.7	1264.0				
0.1704	1279.3	1268.5	1266.9	1269.8	1256.1	1265.0	1265.7	1265.3				
0.2421	1282.4	1273.0	1271.0	1273.4	1262.5	1268.6	1269.7	1268.5				
0.3228	1285.7	1278.2	1276.1	1278.0	1269.6	1273.5	1274.8	1273.1				
0.4163	1289.2	1283.6	1281.6	1283.0	1276.8	1279.1	1280.4	1278.5				
0.5288	1292.0	1288.9	1287.2	1288.1	1283.9	1285.1	1286.2	1284.5				
0.6577	1294.5	1293.5	1292.2	1292.7	1290.2	1290.7	1291.5	1290.1				
0.8089	1296.7	1297.1	1296.4	1296.6	1295.5	1295.6	1296.1	1295.2				
1.0000	1299.8	1299.7	1299.7	1299.7	1299.7	1299.7	1299.7	1299.7				
χ^2		0.3798	0.5288	0.3174	2.5121	0.7545	0.6518	0.7578				

SPT for mixtures of hard convex (not necessarily spherical) molecules gives the equation for mixture as follows [50,51]:

$$
\frac{p}{\rho_N k_B T} = \frac{1}{(1 - V \rho_N)} + \frac{AB \rho_N}{(1 - V \rho_N)^2} + \frac{B^2 C \rho_N^2}{3(1 - V \rho_N)^3}
$$
(8)

where

$$
A = \sum x_i \overline{R_i}, \quad B = \sum x_i S_i, \quad C = \sum x_i \overline{R_i}^2, \quad V = \sum x_i V_{H_i} \tag{9}
$$

 \overline{R}_i , S_i and V_{H_i} are mean radius of curvature, surface area and volume, respectively, of a molecule of species *i*, ρ_N is number density of mixture molecule, and x_i is the mole fraction.

Relating this with the equation:

$$
\gamma \left(\frac{dp}{d\rho}\right)_T = u^2 \tag{10}
$$

where u is the ultrasonic speed, ρ is molecular density and γ is the ratio of specific heats, we get:

$$
\frac{Mu^2}{\gamma RT} = \frac{1}{(1 - V\rho_N)^2} + 2AB \frac{\rho_N}{(1 - V\rho_N)^3} + B^2 C \frac{\rho_N^2}{(1 - V\rho_N)^4}
$$
(11)

Eq. (11) is used to evaluate the ultrasonic speeds in the binary mixtures.

For the case of pure liquids, the above equation is modified by introducing the dimensionless shape parameter, $X = \overline{R}S/V_H$ and $\alpha = V_H \rho_N$:

$$
\frac{Mu^2}{\gamma RT} = \frac{[1 + (X - 1)\alpha]^2}{(1 - \alpha)^4}
$$
\n(12)

Solution to the above equation is obtained as [51]:

$$
\alpha = K - \sqrt{K^2 + L - 1} \tag{13}
$$

where K=1+L(X – 1)/2 and $L = \sqrt{\gamma RT/M u^2}$. The evaluated α values are given in Table 6. Mean r[adius](#page-8-0) and the surface area of a molecule can be written as

$$
\overline{R} = Y V_H^{1/3} \text{ and } S = Z \overline{R}^2 \tag{20}
$$

where Y and Z are the parameters related to the shape of the molecule. If the molecule is assigned different shapes (see Table 4), then the corresponding values of X , Y and Z (shape parameters) can be calculated (see Table 5).

In the present investigation, the theoretical speeds based on SPT is computed for the toluene–nitrile binary mixtures by considering different shapes, viz., sphere, cube, tetrahedro[n,](#page-5-0) [disc](#page-5-0) [A,](#page-5-0) disc B, disc C and disc D for both the participating components. Hence for a binary [mixture](#page-5-0) [t](#page-5-0)here are 49 combinations of shapes at a particular temperature. By using χ^2 test, the best-fit combination of shapes of the participating component of molecules is arrived at. Chi- square

Table 8

Ultrasonic speeds computed theoretically using SPT with behavioural shape having least Chi-square fit and experimentally measured values of the binary mixture TOL + AN at 303.15 and 308.15 K.

x_1	$TOL+AN$							
	303.15K		308.15K					
	u_{exp}	u_{cal} $Disc D + tetra$	u_{exp}	u_{cal} $Disc D + tetra$				
0.0000	1251.9	1251.8	1231.7	1231.6				
0.0506	1253.9	1248.0	1233.4	1227.7				
0.1082	1256.3	1247.4	1235.5	1227.0				
0.1704	1259.0	1249.1	1238.0	1228.7				
0.2421	1261.8	1252.4	1240.9	1232.1				
0.3228	1265.0	1256.7	1244.3	1236.4				
0.4163	1268.2	1261.5	1247.8	1241.4				
0.5288	1270.9	1266.6	1250.7	1246.7				
0.6577	1273.6	1271.4	1253.5	1251.8				
0.8089	1275.9	1275.7	1256.3	1256.5				
1.0000	1279.5	1279.4	1260.7	1260.6				
χ^2		0.34954		0.31716				

Table 9

Ultrasonic speeds computed theoretically using SPT with behavioural shape having least Chi-square fit and experimentally measured values of the binary mixture TOL + PN at 298.15, 303.15 and 308.15 K.

x_1	$TOL + PN$											
	298.15K		303.15K		308.15K							
	u_{\exp}	u_{cal} Disc D + tetra	$u_{\rm exp}$	u_{cal} Disc D + tetra	$u_{\rm exp}$	u_{cal} Disc D + tetra						
0.0000	1252.4	1252.3	1232.7	1232.6	1213.7	1213.6						
0.0668	1257.2	1256.8	1237.0	1236.6	1217.6	1217.3						
0.1397	1262.2	1261.8	1241.7	1241.3	1222.0	1221.8						
0.2164	1267.3	1267.0	1246.5	1246.2	1226.7	1226.7						
0.3011	1272.9	1272.4	1252.0	1251.5	1232.0	1232.0						
0.3905	1278.2	1277.6	1257.1	1256.7	1237.1	1237.3						
0.4917	1283.0	1282.9	1262.1	1262.1	1242.1	1242.7						
0.5989	1287.1	1287.8	1266.2	1267.0	1246.3	1247.8						
0.7174	1291.2	1292.3	1270.4	1271.6	1250.5	1252.6						
0.8540	1295.2	1296.4	1274.5	1275.9	1255.2	1257.0						
1.0000	1299.8	1299.7	1279.5	1279.4	1260.7	1260.6						
χ^2		0.0034		0.0040		0.0082						

Table 10

Ultrasonic speeds computed theoretically using SPT with behavioural shape having least Chi-square fit and experimentally measured values of the binary mixture TOL + AN at 298.15, 303.15 and 308.15 K.

test (χ^2) can be used to test the goodness of fit, which enables us to find whether the deviations of the theoretical values from the experimental ones are due to chance or really due to the inadequacy of the theory to fit the data.

The theoretical ultrasonic speeds based on SPT for the mixtures having liquids belonging to nitrile series with toluene at different temperatures are given in Tables 7–10, for the minimum values of χ^2 . For example, at 298.15 K, for toluene + AN binary mixture 7×7 combinations of the different molecular shapes have been tried for the participating components and only when toluene bears the disc D shape the χ^2 value was found to be least. Hence as in Table 7, we have [given](#page-6-0) [the](#page-6-0) [dis](#page-6-0)c D shape for toluene and oth[er](#page-6-0) [pos](#page-6-0)sible shapes such as sphere, cube, tetrahedron, disc A, disc B, disc C and disc D for acetonitrile. From the χ^2 values (see Table 7), it is clear that when toluene + AN binary mixture takes the shapes of disc D + tetrahedron, the ultrasonic velocity estimated based on SPT is found to be very close to the experimental values.

In order to see whether changes in temperature affects the shapes of the molecules in the mixture, the theoretical ultrasonic speeds of this mixture is evaluated at 303.15 and 308.15 K, respectively, for all the combination of the shape of the molecules (see Table 8), it is found that increase of temperature has no effect on the shape of the participating molecules.

Similarly for the toluene + PN binary mixture (Table 9), the theoretical ultrasonic speeds closely agrees well with the experimental values, when the participating molecules takes the shapes of disc D + tetrahedron at the studied temperatures.

From Table 10, it is further seen that, for the toluene + BN binary mixture, the participating molecules takes the shapes of disc D + disc C. Thus our theoretical analysis using SPT has proved that

the temperature has no effect on the shapes of the participating components in all the three binary systems.

Using the volume interaction parameters (see Table 11) one can conclude that, in toluene + AN/PN binary mixtures, the nitrile molecules expand more than the toluene molecules while vice versa is noted in toluene + BN binary mixture.

Similar conclusions were drawn by Kalidoss et al. [51], using volume interaction parameter[s,](#page-7-0) [for](#page-7-0) [the](#page-7-0) binary mixtures of CCl_4 + benzene, CCl_4 + propanol, CCl_4 + butanol and benzene + butanol binary mixtures.

6. Conclusion

In this paper, density, ultrasonic speed and viscosity data are reported for binary mixtures of toluene and AN/PN/BN at three different temperatures. The excess parameter studies such as Δu , L_f^E , Z^E and ΔG^{*E} indicate the presence of specific donor–acceptor (charge-transfer) interactions between toluene and nitrile molecules. Furthermore, theoretical ultrasonic speeds of these mixtures are evaluated by considering the different shapes of the participating components using the scaled particle theory.

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